

Generalized Wigner-Yanase Skew Information And Generalized Fisher Information

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Abstract

We introduce a generalized Wigner-Yanase skew information and then derive the trace inequality related to the uncertainty relation. This inequality is a non-trivial generalization of the uncertainty relation derived by S.Luo for the quantum uncertainty quantity excluding the classical mixture. And we introduce a generalized Fisher information and then derive a generalized Cramér-Rao inequality. We also give an example for our generalized Fisher information and then derive the uncertainty relation for two observables.

1. INTRODUCTION

As a degree for non-commutativity between a quantum state ρ and an observable H , Wigner-Yanase skew information

$$I_\rho(H) \equiv \frac{1}{2} \text{Tr} \left[\left(i \left[\rho^{1/2}, H \right] \right)^2 \right]$$

was defined in [10]. Here we denote the commutator by $[X, Y] = XY - YX$. This quantity was generalized by Dyson

$$I_{\rho, \alpha}(H) = \frac{1}{2} \text{Tr} \left[(i[\rho^\alpha, H])(i[\rho^{1-\alpha}, H]) \right]$$

which is known as the Wigner-Yanase-Dyson skew information. It is famous that the convexity of $I_{\rho, \alpha}(H)$ with respect to ρ was successfully proven by E.Lieb in [7]. From the physical point of view, an observable H is generally considered to be an unbounded operator, however in the present paper, unless otherwise stated, we consider $H \in B(\mathcal{H})$, where $B(\mathcal{H})$ represents the set of all bounded linear operators on the Hilbert space

\mathcal{H} , as a mathematical interest. We also denote the set of all self-adjoint operators (observables) by $\mathcal{L}_h(\mathcal{H})$ and the set of all density operators (quantum states) by $\mathcal{S}(\mathcal{H})$ on the Hilbert space \mathcal{H} . The relation between the Wigner-Yanase skew information and the uncertainty relation was studied in [9]. Moreover the relation between the Wigner-Yanase-Dyson skew information and the uncertainty relation was studied in [6, 11]. In our previous paper [11], we defined a generalized skew information and then derived a kind of an uncertainty relation. In the section 2, we introduce a new generalized Wigner-Yanase-Dyson skew information. On a generalization of the original Wigner-Yanase-Dyson skew information, our generalization is different from the Wigner-Yanase-Dyson skew information and a generalized skew information defined in our previous paper [11].

On the other hand, we have some definitions for the Fisher information in quantum mechanical system. In the section 3, we consider the standard definition and its one-parameter extended one. For a parameterized density operator $\rho_\theta \in \mathcal{S}_\theta(\mathcal{H})$, we define the Fisher information by

$$I(\rho_\theta, L_\theta) \equiv \text{Tr}[\rho_\theta L_\theta L_\theta^*],$$

where the logarithmic derivative L_θ is defined by

$$\frac{\partial \rho_\theta}{\partial \theta} \equiv \frac{1}{2}(\rho_\theta L_\theta + L_\theta^* \rho_\theta)$$

and $\mathcal{S}_\theta(\mathcal{H})$ represents the set of all quantum states with one-parameter $\theta \in \mathbb{R}$. In the section 3 of the present paper, we define a one-parameter extended Fisher information and study some trace inequalities between this quantity and the variance (a generalized Cramér-Rao type inequality). See the literatures [2, 3] on recent advances of the skew information, the Fisher information and the uncertainty relation.

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2. TRACE INEQUALITIES ON A GENERALIZED WIGNER-YANASE SKEW INFORMATION

We review the relation between the Wigner-Yanase skew information and the uncertainty relation. In quantum mechanical system, the expectation value of an observable H in a quantum state ρ is expressed by $Tr[\rho H]$. It is natural that the variance for a quantum state ρ and an observable H is defined by $V_\rho(H) \equiv Tr[\rho(H - Tr[\rho H]I)^2] = Tr[\rho H^2] - Tr[\rho H]^2$. It is famous that we have the Heisenberg's uncertainty relation:

$$V_\rho(A)V_\rho(B) \geq \frac{1}{4}|Tr[\rho[A, B]]|^2 \quad (1)$$

for a quantum state ρ and two observables A and B . The further strong result was given by Schrödinger

$$V_\rho(A)V_\rho(B) - |Cov_\rho(A, B)|^2 \geq \frac{1}{4}|Tr[\rho[A, B]]|^2,$$

where the covariance is defined by $Cov_\rho(A, B) \equiv Tr[\rho(A - Tr[\rho A]I)(B - Tr[\rho B]I)]$. However, the uncertainty relation for the skew information failed. (See [9, 6, 11].)

$$I_\rho(A)I_\rho(B) \geq \frac{1}{4}|Tr[\rho[A, B]]|^2.$$

Recently S.Luo introduced the quantity $U_\rho(H)$ representing a quantum uncertainty excluding the classical mixture:

$$U_\rho(H) \equiv \sqrt{V_\rho(H)^2 - (V_\rho(H) - I_\rho(H))^2}.$$

Note that we have the relation among quantities as

$$0 \leq I_\rho(H) \leq U_\rho(H) \leq V_\rho(H). \quad (2)$$

For a quantum state ρ and observables X, Y , he derived the following uncertainty relation in [8]:

$$U_\rho(X)U_\rho(Y) \geq \frac{1}{4}|Tr[\rho[X, Y]]|^2. \quad (3)$$

The inequality (3) is a refinement of the inequality (1) in the sense of (2). In this section, we study two types of one-parameter extended inequalities for the inequality (3).

Definition 2.1 For $0 \leq \alpha \leq 1$, a quantum state ρ and an observable H , we define the Wigner-Yanase-Dyson skew information

$$I_{\rho, \alpha}(H) \equiv \frac{1}{2}Tr [(i[\rho^\alpha, H_0])(i[\rho^{1-\alpha}, H_0])]$$

and we also define

$$J_{\rho, \alpha}(H) \equiv \frac{1}{2}Tr [\{\rho^\alpha, H_0\}\{\rho^{1-\alpha}, H_0\}],$$

where $H_0 \equiv H - Tr[\rho H]I$ and we denote the anti-commutator by $\{X, Y\} = XY + YX$.

Note that we have

$$\begin{aligned} & \frac{1}{2}Tr [(i[\rho^\alpha, H_0])(i[\rho^{1-\alpha}, H_0])] \\ &= \frac{1}{2}Tr [(i[\rho^\alpha, H])(i[\rho^{1-\alpha}, H])], \end{aligned}$$

but we have

$$\frac{1}{2}Tr [\{\rho^\alpha, H_0\}\{\rho^{1-\alpha}, H_0\}] \neq \frac{1}{2}Tr [\{\rho^\alpha, H\}\{\rho^{1-\alpha}, H\}].$$

Then we have the following inequalities:

$$I_{\rho, \alpha}(H) \leq I_\rho(H) \leq J_\rho(H) \leq J_{\rho, \alpha}(H), \quad (4)$$

since we have $Tr[\rho^{1/2}H\rho^{1/2}H] \leq Tr[\rho^\alpha H\rho^{1-\alpha}H]$. If we define

$$U_{\rho, \alpha}(H) \equiv \sqrt{V_\rho(H)^2 - (V_\rho(H) - I_{\rho, \alpha}(H))^2},$$

as a direct generalization of Eq.(1), then we have

$$0 \leq I_{\rho, \alpha}(H) \leq U_{\rho, \alpha}(H) \leq U_\rho(H)$$

due to the first inequality of (9). We also have

$$U_{\rho, \alpha}(H) = \sqrt{I_{\rho, \alpha}(H)J_{\rho, \alpha}(H)}.$$

In this paper, we introduce a generalized Wigner-Yanase skew information which is a generalized Wigner-Yanase skew information by

$$K_{\rho, \alpha}(H) \equiv \frac{1}{2}Tr \left[\left(i \left[\frac{\rho^\alpha + \rho^{1-\alpha}}{2}, H_0 \right] \right)^2 \right]$$

and we also define

$$L_{\rho, \alpha}(H) \equiv \frac{1}{2}Tr \left[\left(\left\{ \frac{\rho^\alpha + \rho^{1-\alpha}}{2}, H_0 \right\} \right)^2 \right].$$

Throughout this section, we put $X_0 \equiv X - Tr[\rho X]I$ and $Y_0 \equiv Y - Tr[\rho Y]I$. Then we show the following trace inequality.

Theorem 2.2 For a quantum state ρ and observable X, Y and $\alpha \in [0, 1]$, we have

$$W_{\rho, \alpha}(X)W_{\rho, \alpha}(Y) \geq \frac{1}{4} \left| Tr \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right) [X, Y] \right] \right|^2,$$

where

$$W_{\rho, \alpha}(X) \equiv \sqrt{K_{\rho, \alpha}(X)L_{\rho, \alpha}(X)}.$$

Remark 2.3 *Theorem 2.2 is not trivial by the following two reasons.*

(1) *There is no relation between*

$$\left| \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right|^2$$

$$\text{and } |\text{Tr} [\rho [X, Y]]|^2.$$

(2) *Though $U_{\rho, \alpha}(H) \leq U_\rho(H)$ and $U_{\rho, \alpha}(H) \leq \tilde{U}_{\rho, \alpha}(H)$ hold, there is no relation between $U_\rho(H)$ and $\tilde{U}_{\rho, \alpha}(H)$.*

Proof of Theorem 2.2. We put

$$A_\alpha(H) \equiv i[\rho^\alpha, H_0], B_\alpha(H) \equiv \{\rho^\alpha, H_0\},$$

$$K = \frac{1}{2}(A_\alpha(X) + A_{1-\alpha}(X))x + \frac{1}{2}(B_\alpha(Y) + B_{1-\alpha}(Y)).$$

It follows from $K^* = K$ that

$$\begin{aligned} 0 &\leq \text{Tr}[KK^*] \\ &= \frac{1}{4}\text{Tr}[(A_\alpha(X) + A_{1-\alpha}(X))^2]x^2 \\ &\quad + \frac{1}{2}\text{Tr}[(A_\alpha(X) + A_{1-\alpha}(X))(B_\alpha(Y) + B_{1-\alpha}(Y))]x \\ &\quad + \frac{1}{4}\text{Tr}[(B_\alpha(Y) + B_{1-\alpha}(Y))^2] \\ &= \left(\frac{1}{4}\text{Tr}[A_\alpha(X)^2 + A_{1-\alpha}(X)^2] + I_{\rho, \alpha}(X) \right) x^2 \\ &\quad + \frac{1}{2}\text{Tr}[(A_\alpha(X) + A_{1-\alpha}(X))(B_\alpha(Y) + B_{1-\alpha}(Y))]x \\ &\quad + \left(\frac{1}{4}\text{Tr}[B_\alpha(Y)^2 + B_{1-\alpha}(Y)^2] + J_{\rho, \alpha}(Y) \right). \end{aligned}$$

Then

$$\begin{aligned} &\frac{1}{4}(\text{Tr}[(A_\alpha(X) + A_{1-\alpha}(X))(B_\alpha(Y) + B_{1-\alpha}(Y))])^2 \\ &\leq 4 \left(\frac{1}{4}\text{Tr}[A_\alpha(X)^2 + A_{1-\alpha}(X)^2] + I_{\rho, \alpha}(X) \right) \quad (5) \\ &\quad \left(\frac{1}{4}\text{Tr}[B_\alpha(Y)^2 + B_{1-\alpha}(Y)^2] + J_{\rho, \alpha}(Y) \right). \end{aligned}$$

Now we have

$$\begin{aligned} &\text{Tr}[(A_\alpha(X) + A_{1-\alpha}(X))(B_\alpha(Y) + B_{1-\alpha}(Y))] \\ &= \text{Tr}[i[\rho^\alpha, X_0] + i[\rho^{1-\alpha}, X_0]](\{\rho^\alpha, Y_0\} + \{\rho^{1-\alpha}, Y_0\}) \\ &= i\text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 X_0 Y_0 - Y_0 X_0 (\rho^\alpha + \rho^{1-\alpha})^2] \\ &= \text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X_0, Y_0])] \\ &= \text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X, Y])]. \end{aligned}$$

Then (5) is equivalent to the following;

$$\begin{aligned} &\frac{1}{4}(\text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X, Y])])^2 \quad (6) \\ &\leq 4 \left(\frac{1}{4}\text{Tr}[A_\alpha(X)^2 + A_{1-\alpha}(X)^2] + I_{\rho, \alpha}(X) \right) \\ &\quad \left(\frac{1}{4}\text{Tr}[B_\alpha(Y)^2 + B_{1-\alpha}(Y)^2] + J_{\rho, \alpha}(Y) \right). \end{aligned}$$

And we also have

$$\begin{aligned} &\frac{1}{4}|\text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X, Y])]|^2 \quad (7) \\ &\leq 4 \left(\frac{1}{4}\text{Tr}[A_\alpha(Y)^2 + A_{1-\alpha}(Y)^2] + I_{\rho, \alpha}(Y) \right) \\ &\quad \left(\frac{1}{4}\text{Tr}[B_\alpha(X)^2 + B_{1-\alpha}(X)^2] + J_{\rho, \alpha}(X) \right). \end{aligned}$$

By taking a square root of (6) \times (7), we have

$$\begin{aligned} &\left\{ \frac{1}{4}(\text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X, Y])])^2 \right\}^2 \\ &\leq 4 \left(\frac{1}{4}\text{Tr}[A_\alpha(X)^2 + A_{1-\alpha}(X)^2] + I_{\rho, \alpha}(X) \right) \\ &\quad \left(\frac{1}{4}\text{Tr}[B_\alpha(Y)^2 + B_{1-\alpha}(Y)^2] + J_{\rho, \alpha}(Y) \right) \\ &\quad 4 \left(\frac{1}{4}\text{Tr}[A_\alpha(Y)^2 + A_{1-\alpha}(Y)^2] + I_{\rho, \alpha}(Y) \right) \\ &\quad \left(\frac{1}{4}\text{Tr}[B_\alpha(X)^2 + B_{1-\alpha}(X)^2] + J_{\rho, \alpha}(X) \right). \end{aligned}$$

Thus

$$\begin{aligned} &\frac{1}{4}(\text{Tr}[(\rho^\alpha + \rho^{1-\alpha})^2 (i[X, Y])])^2 \\ &\leq 2\sqrt{\left(\frac{1}{4}\text{Tr}[A_\alpha(X)^2 + A_{1-\alpha}(X)^2] + I_{\rho, \alpha}(X) \right)} \\ &\quad \sqrt{\left(\frac{1}{4}\text{Tr}[B_\alpha(Y)^2 + B_{1-\alpha}(Y)^2] + J_{\rho, \alpha}(Y) \right)} \\ &\quad 2\sqrt{\left(\frac{1}{4}\text{Tr}[A_\alpha(Y)^2 + A_{1-\alpha}(Y)^2] + I_{\rho, \alpha}(Y) \right)} \\ &\quad \sqrt{\left(\frac{1}{4}\text{Tr}[B_\alpha(X)^2 + B_{1-\alpha}(X)^2] + J_{\rho, \alpha}(X) \right)}. \end{aligned}$$

Therefore

$$\begin{aligned} &\frac{1}{4} \left(\text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 (i[X, Y]) \right] \right)^2 \\ &\leq W_\alpha(\rho, X)W_\alpha(\rho, Y). \end{aligned}$$

Since

$$\begin{aligned} & \overline{\text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right]} \\ &= -\text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right], \end{aligned}$$

we have

$$\text{Re} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] = 0.$$

And then

$$\begin{aligned} & \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \\ &= i \text{ImTr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right]. \end{aligned}$$

Hence

$$\begin{aligned} & \left(\text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 (i[X, Y]) \right] \right)^2 \\ &= - \left(\text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right)^2 \\ &= - \left(i \text{ImTr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right)^2 \\ &= \left(\text{ImTr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right)^2 \\ &= \left| \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right|^2. \end{aligned}$$

q.e.d.

We also define the followings to obtain another uncertainty relation.

Definition 2.4 For a quantum state ρ and observable H and $\alpha \in [0, 1]$, we define

$$\begin{aligned} & \tilde{W}_{\rho, \alpha}(H) \\ &\equiv \frac{1}{4} \sqrt{\text{Tr} [(i[\rho^\alpha, H_0])^2] \text{Tr} [(i[\rho^{1-\alpha}, H_0])^2]} \\ & \quad \sqrt{\text{Tr} [\{\rho^\alpha, H_0\}^2] \text{Tr} [\{\rho^{1-\alpha}, H_0\}^2]}. \end{aligned}$$

The we have the following theorem.

Theorem 2.5 For a quantum state ρ and observable X, Y and $\alpha \in [0, 1]$, we have

$$\begin{aligned} & \sqrt{\tilde{W}_{\rho, \alpha}(X) \tilde{W}_{\rho, \alpha}(Y)} \\ & \geq \frac{1}{4} \left| \text{Tr} [\rho^{2\alpha} [X, Y]] \text{Tr} [\rho^{2(1-\alpha)} [X, Y]] \right|. \end{aligned}$$

Remark 2.6 There is no relation between Theorem 2.2 and Theorem 2.5 by the following (1), (2).

(1) There is no relation between $4\tilde{W}_{\rho, \alpha}(X)$ and

$$\begin{aligned} & \left(\text{Tr} \left[\frac{(i[\rho^\alpha, X_0])^2 + (i[\rho^{1-\alpha}, X_0])^2}{4} \right] + I_{\rho, \alpha}(X) \right) \\ & \left(\text{Tr} \left[\frac{(\{\rho^\alpha, X_0\})^2 + (\{\rho^{1-\alpha}, X_0\})^2}{4} \right] + J_{\rho, \alpha}(X) \right). \end{aligned}$$

That is, there are no relation between

$$\sqrt{\text{Tr} [(i[\rho^\alpha, X_0])^2] \text{Tr} [(i[\rho^{1-\alpha}, X_0])^2]}$$

and

$$\begin{aligned} & \text{Tr} \left[\frac{(i[\rho^\alpha, X_0])^2 + (i[\rho^{1-\alpha}, X_0])^2}{4} \right] \\ & + \frac{1}{2} \text{Tr} [(i[\rho^\alpha, X_0])(i[\rho^{1-\alpha}, X_0])]. \end{aligned}$$

and there is no relation between

$$\sqrt{\text{Tr} [\{\rho^\alpha, X_0\}^2] \text{Tr} [\{\rho^{1-\alpha}, X_0\}^2]}$$

and

$$\begin{aligned} & \text{Tr} \left[\frac{\{\rho^\alpha, X_0\}^2 + \{\rho^{1-\alpha}, X_0\}^2}{4} \right] \\ & + \frac{1}{2} \text{Tr} [\{\rho^\alpha, X_0\} \{\rho^{1-\alpha}, X_0\}]. \end{aligned}$$

(2) There is no relation between

$$\left| \text{Tr} [\rho^{2\alpha} [X, Y]] \text{Tr} [\rho^{2(1-\alpha)} [X, Y]] \right|$$

and

$$\left| \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right|^2.$$

That is, there is no relation between

$$\left| \text{Tr} [\rho^{2\alpha} [X, Y]] \right|$$

and

$$\left| \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right|.$$

and there in no relation between

$$\left| \text{Tr} \left[\rho^{2(1-\alpha)} [X, Y] \right] \right|$$

and

$$\left| \text{Tr} \left[\left(\frac{\rho^\alpha + \rho^{1-\alpha}}{2} \right)^2 [X, Y] \right] \right|.$$

(3) When $\alpha = 1/2$, both Theorem 2.2 and Theorem 2.5 reduce the result of Luo.

Proof of Theorem 2.5. We put

$$K = i[\rho^\alpha, X_0]x + \{\rho^\alpha, Y_0\}.$$

It follows from $K^* = K$ that

$$\begin{aligned} 0 &\leq \text{Tr} [KK^*] \\ &= \text{Tr} [(i[\rho^\alpha, X_0]x + \{\rho^\alpha, Y_0\})^2] \\ &= \text{Tr} [(i[\rho^\alpha, X_0])^2] x^2 + 2i\text{Tr} [[\rho^\alpha, X_0]\{\rho^\alpha, Y_0\}] x \\ &\quad + \text{Tr} [\{\rho^\alpha, Y_0\}^2] \\ &= \text{Tr} [(i[\rho^\alpha, X_0])^2] x^2 + 2i\text{ImTr} [\rho^{2\alpha} [X, Y]] x \\ &\quad + \text{Tr} [\{\rho^\alpha, Y_0\}^2]. \end{aligned}$$

Then

$$\begin{aligned} |\text{Tr} [\rho^{2\alpha} [X, Y]]|^2 &= (\text{ImTr} [\rho^{2\alpha} [X, Y]])^2 \\ &\leq \text{Tr} [(i[\rho^\alpha, X_0])^2] \text{Tr} [\{\rho^\alpha, Y_0\}^2]. \end{aligned}$$

By exchanging X and Y we have

$$\begin{aligned} |\text{Tr} [\rho^{2\alpha} [X, Y]]|^2 &\leq \text{Tr} [(i[\rho^\alpha, Y_0])^2] \text{Tr} [\{\rho^\alpha, X_0\}^2]. \end{aligned}$$

And we also have

$$\begin{aligned} |\text{Tr} [\rho^{2(1-\alpha)} [X, Y]]|^2 &\leq \text{Tr} [(i[\rho^{1-\alpha}, X_0])^2] \text{Tr} [\{\rho^{1-\alpha}, Y_0\}^2]. \end{aligned}$$

By exchanging X and Y we have

$$\begin{aligned} |\text{Tr} [\rho^{2(1-\alpha)} [X, Y]]|^2 &\leq \text{Tr} [(i[\rho^{1-\alpha}, Y_0])^2] \text{Tr} [\{\rho^{1-\alpha}, X_0\}^2]. \end{aligned}$$

We put as follows;

$$S_{\rho, \alpha}(X) \equiv \frac{1}{2} \text{Tr} [(i[\rho^\alpha, X_0])^2],$$

$$S_{\rho, 1-\alpha}(X) \equiv \frac{1}{2} \text{Tr} [(i[\rho^{1-\alpha}, X_0])^2],$$

$$S_{\rho, \alpha}(Y) \equiv \frac{1}{2} \text{Tr} [(i[\rho^\alpha, Y_0])^2],$$

$$S_{\rho, 1-\alpha}(Y) \equiv \frac{1}{2} \text{Tr} [(i[\rho^{1-\alpha}, Y_0])^2],$$

$$T_{\rho, \alpha}(X) \equiv \frac{1}{2} \text{Tr} [\{\rho^\alpha, X_0\}^2],$$

$$T_{\rho, 1-\alpha}(X) \equiv \frac{1}{2} \text{Tr} [\{\rho^{1-\alpha}, X_0\}^2],$$

$$T_{\rho, \alpha}(Y) \equiv \frac{1}{2} \text{Tr} [\{\rho^\alpha, Y_0\}^2],$$

$$T_{\rho, 1-\alpha}(Y) \equiv \frac{1}{2} \text{Tr} [\{\rho^{1-\alpha}, Y_0\}^2].$$

Then we have

$$\begin{aligned} &|\text{Tr} [\rho^{2\alpha} [X, Y]]|^2 \\ &\leq 4\sqrt{S_{\rho, \alpha}(X)T_{\rho, \alpha}(X)S_{\rho, \alpha}(Y)T_{\rho, \alpha}(Y)}. \end{aligned}$$

$$|\text{Tr} [\rho^{2(1-\alpha)} [X, Y]]|^2$$

$$\leq 4\sqrt{S_{\rho, 1-\alpha}(X)T_{\rho, 1-\alpha}(X)S_{\rho, 1-\alpha}(Y)T_{\rho, 1-\alpha}(Y)}.$$

By putting

$$\tilde{W}_{\rho, \alpha}(X) \equiv \sqrt{S_{\rho, \alpha}(X)S_{\rho, 1-\alpha}(X)T_{\rho, \alpha}(X)T_{\rho, 1-\alpha}(X)},$$

$$\tilde{W}_{\rho, \alpha}(Y) \equiv \sqrt{S_{\rho, \alpha}(Y)S_{\rho, 1-\alpha}(Y)T_{\rho, \alpha}(Y)T_{\rho, 1-\alpha}(Y)},$$

we have

$$\begin{aligned} &\sqrt{\tilde{W}_{\rho, \alpha}(X)\tilde{W}_{\rho, \alpha}(Y)} \\ &\geq \frac{1}{4} |\text{Tr} [\rho^{2\alpha} [X, Y]] \text{Tr} [\rho^{2(1-\alpha)} [X, Y]]|. \end{aligned}$$

q.e.d.

3. A GENERALIZED FISHER INFORMATION AND A GENERALIZED CRAMÉR-RAO INEQUALITY

We review the Fisher information and the Cramér-Rao inequality in quantum mechanical system. We consider the set of all quantum states:

$$\mathcal{S}_\theta(\mathcal{H}) \equiv \{\rho_\theta \in B(\mathcal{H}) | \rho_\theta \geq 0, \text{Tr}[\rho_\theta] = 1\},$$

with one parameter $\theta \in \mathbb{R}$. Let $H \in \mathcal{L}_h(\mathcal{H}) \equiv \{H \in B(\mathcal{H}) | H = H^*\}$ be an estimator of the parameter θ . In the sequel, we consider the case which an estimator is unbiased, that is, $E_\theta[H] \equiv \text{Tr}[\rho_\theta H] = \theta$. The variance $V_\theta[H]$ of the estimator H is defined by $V_\theta[H] \equiv \text{Tr}[\rho_\theta(H - \text{Tr}[\rho_\theta H])^2]$. Then the famous Cramér-Rao inequality, which is a relation between the Fisher information and the variance, $V_\theta[H] \geq \frac{1}{I(\rho_\theta, L_\theta)}$ holds. We should note that the logarithmic derivative

$L_\theta \in B(\mathcal{H})$ is not uniquely determined. Thus we define the symmetric logarithmic derivative $L_\theta^S \in \mathcal{L}_h(\mathcal{H})$ by

$$\frac{\partial \rho_\theta}{\partial \theta} \equiv \frac{1}{2}(\rho_\theta L_\theta^S + L_\theta^S \rho_\theta).$$

Then the symmetric logarithmic derivative L_θ^S is uniquely determined [1, 4, 5] and we have

$$I(\rho_\theta, L_\theta) \geq I(\rho_\theta, L_\theta^S). \quad (8)$$

In addition, for the symmetric logarithmic derivative L_θ^S , we have the Cramér-Rao inequality [1, 4, 5]:

$$V_\theta[H] \geq \frac{1}{I(\rho_\theta, L_\theta^S)}. \quad (9)$$

Due to the inequality (8), we have the following theorem known as Cramér-Rao inequality.

Theorem 3.1

$$V_\theta[H] \geq \frac{1}{I(\rho_\theta, L_\theta)}.$$

That is, the symmetric logarithmic derivative L_θ^S gives the best estimation of the lower bound for the variance $V_\theta[H]$.

We here introduce a generalized Fisher information with one-parameter $\alpha \in [0, 1]$.

Definition 3.2 We define a generalized Fisher information by

$$I_\alpha(\rho_\theta, L_{\theta, \alpha}) \equiv \text{Tr}[\rho_\theta^\alpha L_{\theta, \alpha} \rho_\theta^{1-\alpha} L_{\theta, \alpha}^*], \quad \alpha \in [0, 1],$$

where a generalized logarithmic derivative $L_{\theta, \alpha}$ is defined by

$$\frac{\partial \rho_\theta}{\partial \theta} \equiv \frac{1}{2} \left(\rho_\theta^{\frac{1+\alpha}{2}} L_{\theta, \alpha} \rho_\theta^{\frac{1-\alpha}{2}} + \rho_\theta^{\frac{1-\alpha}{2}} L_{\theta, \alpha}^* \rho_\theta^{\frac{1+\alpha}{2}} \right). \quad (10)$$

Note that $\alpha = 1$ or $[\rho_\theta, L_{\theta, \alpha}] = 0$ recovers $I_\alpha(\rho_\theta, L_{\theta, \alpha}) = I(\rho_\theta, L_\theta)$. We also have $I_\alpha(\rho_\theta, L_{\theta, \alpha}) \geq 0$ and the following trace inequality.

Theorem 3.3 For a self-adjoint operator H , a density operator ρ_θ with the parameter θ and $\alpha \in [0, 1]$, if we have $E_\theta[H] = \theta$, then we have the inequality

$$V_\theta[H] \geq \frac{1}{I_\alpha(\rho_\theta, L_{\theta, \alpha})}. \quad (11)$$

It is clear that (11) is obtained by putting

$$L_{\rho, \theta} = \rho_\theta^{\frac{1-\alpha}{2}} L_{\theta, \alpha} \rho_\theta^{\frac{\alpha-1}{2}}.$$

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