Generalized Wigner-Yanase-Dyson Skew Information and Uncertainty Relation

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Abstract—We give a trace inequality related to the uncertainty relation of generalized Wigner-Yanase-Dyson skew information which includes our result in [16].

I. INTRODUCTION

It is known that the relation between quantum covariances and quantum Fisher informations is studied and the study is applied to generalize a recently proved uncertainty relation based on quantum Fisher information. For example see [1], [6], [7]. Wigner-Yanase skew information

$$I_{\rho}(H) = \frac{1}{2}Tr\left[\left(i\left[\rho^{1/2}, H\right]\right)^{2}\right]$$

= $Tr[\rho H^{2}] - Tr[\rho^{1/2}H\rho^{1/2}H]$

was defined in [13]. This quantity can be considered as a kind of the degree for non-commutativity between a quantum state ρ and an observable H. Here we denote the commutator by [X, Y] = XY - YX. This quantity was generalized by Dyson to

$$I_{\rho,\alpha}(H) = \frac{1}{2}Tr[(i[\rho^{\alpha}, H])(i[\rho^{1-\alpha}, H])]$$

= $Tr[\rho H^2] - Tr[\rho^{\alpha} H \rho^{1-\alpha} H], \alpha \in [0, 1]$

which is known as the Wigner-Yanase-Dyson skew information. It is famous that the convexity of $I_{\rho,\alpha}(H)$ with respect to ρ was successfully proven by E.H.Lieb in [10]. And also this quantity was generalized by Chen and Luo in [4] to

$$I_{\rho,\alpha,\beta}(H)$$

$$= \frac{1}{2}Tr[(i[\rho^{\alpha},H])(i[\rho^{\beta},H])\rho^{1-\alpha-\beta}]$$

$$= \frac{1}{2}\{Tr[\rho H^{2}] + Tr[\rho^{\alpha+\beta}H\rho^{1-\alpha-\beta}H]$$

$$-Tr[\rho^{\alpha}H\rho^{1-\alpha}H] - Tr[\rho^{\beta}H\rho^{1-\beta}H]\},$$

where $\alpha, \beta \geq 0, \alpha + \beta \leq 1$. The convexity of $I_{\rho,\alpha,\beta}(H)$ with respect to ρ was proven by Cai and Luo in [3] under some restrictive condition. From the physical point of view, an observable H is generally considered to be an unbounded opetrator, however in the present paper, unless otherwise stated, we consider $H \in B(\mathcal{H})$ (the set of all bounded linear operators on the Hilbert space \mathcal{H}) as a mathematical interest. We also denote the set of all self-adjoint operators (observables) by $\mathcal{L}_h(\mathcal{H})$ and the set of all density operators (quantum states) by $\mathcal{S}(\mathcal{H})$ on the Hilbert space \mathcal{H} . The relation between the Wigner-Yanase skew information and the uncertainty relation was studied in [12]. Moreover the relation between the Wigner-Yanase-Dyson skew information and the uncertainty relation was studied in [9], [14]. In our paper [14] and [16], we defined a generalized skew information and then derived a kind of uncertainty relation. In section 2, we discuss various properties of Wigner-Yanase-Dyson skew information. In section 3, we give an uncertainty relation of generalized Wigner-Yanase-Dyson skew information.

II. TRACE INEQUALITY OF WIGNER-YANASE-DYSON SKEW INFORMATION

We review the relation between the Wigner-Yanase skew information and the uncertainty relation. In quantum mechanical system, the expectation value of an observable H in a quantum state ρ is expressed by $Tr[\rho H]$. It is natural that the variance for a quantum state ρ and an observable H is defined by $V_{\rho}(H) = Tr[\rho(H - Tr[\rho H]I)^2] = Tr[\rho H^2] - Tr[\rho H]^2$. It is famous that we have

$$V_{\rho}(A)V_{\rho}(B) \ge \frac{1}{4}|Tr[\rho[A,B]]|^2$$
 (1)

for a quantum state ρ and two observables A and B. The further strong results was given by Schrödinger

$$V_{\rho}(A)V_{\rho}(B) - |Cov_{\rho}(A,B)|^{2} \ge \frac{1}{4}|Tr[\rho[A,B]]|^{2},$$

where the covariance is defined by $Cov_{\rho}(A, B) = Tr[\rho(A - Tr[\rho A]I)(B - Tr[\rho B]I)]$. However, the uncertainty relation for the Wigner-Yanase skew information failed. (See [12], [9], [14])

$$I_{\rho}(A)I_{\rho}(B) \ge \frac{1}{4}|Tr[\rho[A,B]]|^{2}.$$

Recently, S.Luo introduced the quantity $U_{\rho}(H)$ representing a quantum uncertainty excluding the classical mixture:

$$U_{\rho}(H) = \sqrt{V_{\rho}(H)^2 - (V_{\rho}(H) - I_{\rho}(H))^2},$$
 (2)

then he derived the uncertainty relation on $U_{\rho}(H)$ in [11]:

$$U_{\rho}(A)U_{\rho}(B) \ge \frac{1}{4}|Tr[\rho[A,B]]|^{2}.$$
 (3)

Note that we have the following relation

$$0 \le I_{\rho}(H) \le U_{\rho}(H) \le V_{\rho}(H). \tag{4}$$

The inequality (3) is a refinement of the inequality (1) in the sense of (4). In [16], we studied one-parameter extended inequality for the inequality (3).

Definition 2.1: For $0 \le \alpha \le 1$, a quantum state ρ and an observable H, we define the Wigner-Yanase-Dyson skew information

$$I_{\rho,\alpha}(H) = \frac{1}{2} Tr[(i[\rho^{\alpha}, H_0])(i[\rho^{1-\alpha}, H_0])]$$

= $Tr[\rho H_0^2] - Tr[\rho^{\alpha} H_0 \rho^{1-\alpha} H_0]$ (5)

and we also define

$$J_{\rho,\alpha}(H) = \frac{1}{2} Tr[\{\rho^{\alpha}, H_0\}\{\rho^{1-\alpha}, H_0\}] = Tr[\rho H_0^2] + Tr[\rho^{\alpha} H_0 \rho^{1-\alpha} H_0], \quad (6)$$

where $H_0 = H - Tr[\rho H]I$ and we denote the anti-commutator by $\{X, Y\} = XY + YX$.

Note that we have

$$\frac{1}{2}Tr[(i[\rho^{\alpha}, H_0])(i[\rho^{1-\alpha}, H_0])] = \frac{1}{2}Tr[(i[\rho^{\alpha}, H])(i[\rho^{1-\alpha}, H])]$$
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but we have

$$\frac{1}{2}Tr[\{\rho^{\alpha}, H_0\}\{\rho^{1-\alpha}, H_0\}] \neq \frac{1}{2}Tr[\{\rho^{\alpha}, H\}\{\rho^{1-\alpha}, H\}].$$

Then we have the following inequalities:

$$I_{\rho,\alpha}(H) \le I_{\rho}(H) \le J_{\rho}(H) \le J_{\rho,\alpha}(H), \tag{7}$$

since we have $Tr[\rho^{1/2}H\rho^{1/2}H] \leq Tr[\rho^{\alpha}H\rho^{1-\alpha}H]$. (See [2], [5] for example.) If we define

$$U_{\rho,\alpha}(H) = \sqrt{V_{\rho}(H)^2 - (V_{\rho}(H) - I_{\rho,\alpha}(H))^2},$$
 (8)

as a direct generalization of Eq.(2), then we have

$$0 \le I_{\rho,\alpha}(H) \le U_{\rho,\alpha}(H) \le U_{\rho}(H) \tag{9}$$

due to the first inequality of (7). We also have

$$U_{\rho,\alpha}(H) = \sqrt{I_{\rho,\alpha}(H)J_{\rho,\alpha}(H)}.$$

From the inequalities (4),(8),(9), our situation is that we have

$$0 \le I_{\rho,\alpha}(H) \le I_{\rho}(H) \le U_{\rho}(H)$$

and

$$0 \le I_{\rho,\alpha}(H) \le U_{\rho,\alpha}(H) \le U_{\rho}(H).$$

We gave the following uncertainty relation with respect to $U_{\rho,\alpha}(H)$ as a direct generalization of the inequality (3).

Theorem 2.1 ([16]): For $0 \le \alpha \le 1$, a quantum state ρ and observablea A, B,

$$U_{\rho,\alpha}(A)U_{\rho,\alpha}(B) \ge \alpha(1-\alpha)|Tr[\rho[A,B]]|^2.$$
(10)

Now we define the two parameter extensions of Wigner-Yanase skew information and give an uncertainty relation under some conditions in the next section.

Definition 2.2: For $\alpha, \beta \geq 0$, a quantum state ρ and an observable H, we define the generalized Wigner-Yanase-Dyson skew information

$$I_{\rho,\alpha,\beta}(H) = \frac{1}{2} Tr \left[(i[\rho^{\alpha}, H_0])(i[\rho^{\beta}, H_0])\rho^{1-\alpha-\beta} \right] \\ = \frac{1}{2} \{ Tr[\rho H_0^2] + Tr[\rho^{\alpha+\beta} H_0 \rho^{1-\alpha-\beta} H_0] \\ - Tr[\rho^{\alpha} H_0 \rho^{1-\alpha} H_0] - Tr[\rho^{\beta} H_0 \rho^{1-\beta} H_0] \}$$

and we define

$$\begin{aligned} &J_{\rho,\alpha,\beta}(H) \\ &= \frac{1}{2} Tr\left[(i\{\rho^{\alpha},H_0\})(i\{\rho^{\beta},H_0\})\rho^{1-\alpha-\beta}\right] \\ &= \frac{1}{2} \{Tr[\rho H_0^2] + Tr[\rho^{\alpha+\beta}H_0\rho^{1-\alpha-\beta}H_0] \\ &+ Tr[\rho^{\alpha}H_0\rho^{1-\alpha}H_0] + Tr[\rho^{\beta}H_0\rho^{1-\beta}H_0]\}, \end{aligned}$$

where $H_0 = H - Tr[\rho H]I$ and we denote the anti-commutator by $\{X, Y\} = XY + YX$. We remark that $\alpha + \beta = 1$ implies $I_{\rho,\alpha}(H) = I_{\rho,\alpha,1-\alpha}(H)$ and $J_{\rho,\alpha}(H) = J_{\rho,\alpha,1-\alpha}(H)$. We also define

$$U_{\rho,\alpha,\beta}(H) = \sqrt{I_{\rho,\alpha,\beta}(H)J_{\rho,\alpha,\beta}(H)}.$$

III. MAIN THEOREM

In this section we assume that ρ is an invertible density matrix and A, B are Hermitian matrices. We also assume that $\alpha, \beta \geq 0$ do not necessarily satisfy the condition $\alpha + \beta \leq 1$. We give the main theorem as follows;

Theorem 3.1: For
$$\alpha, \beta \ge 0$$
 and $\alpha + \beta \ge 1$ or $\alpha + \beta \le \frac{1}{2}$,
 $U_{\rho,\alpha,\beta}(A)U_{\rho,\alpha,\beta}(B) \ge \alpha\beta |Tr[\rho[A,B]]|^2$. (11)

In order to prove Theorem 3.1, we use the several lemmas. By spectral decomposition, there exists an orthonormal basis $\{\phi_1, \phi_2, \dots, \phi_n\}$ consisting of eigenvectors of ρ . Let $\lambda_1, \lambda_2, \ldots, \lambda_n$ be the corresponding eigenvalues, where $\sum_{i=1}^{n} \lambda_i = 1$ and $\lambda_i > 0$. Thus, ρ has a spectral representation

$$\rho = \sum_{i=1}^{n} \lambda_i |\phi_i\rangle \langle \phi_i|.$$
(12)

We use the notation $f_{\alpha}(i,j) = \lambda_i^{\alpha} \lambda_j^{1-\alpha} + \lambda_i^{1-\alpha} \lambda_j^{\alpha}$. And we also use $h_{ij} = \langle \phi_i | H_0 | \phi_j \rangle$, $a_{ij} = \langle \phi_i | A_0 | \phi_j \rangle$ and $b_{ij} =$ $\langle \phi_i | B_0 | \phi_i \rangle$. Then we have the following lemmas. *Lemma 3.1:*

$$= \frac{I_{\rho,\alpha,\beta}(H)}{2} \sum_{i < j} \{\lambda_i + \lambda_j + f_{\alpha+\beta}(i,j) - f_{\alpha}(i,j) - f_{\beta}(i,j)\} |h_{ij}|^2$$

Proof of Lemma 3.1. By (12),

$$\rho H_0^2 = \sum_{i=1}^n \lambda_i |\phi_i\rangle \langle \phi_i | H_0^2.$$

Then

$$Tr[\rho H_0^2] = \sum_{i=1}^n \lambda_i \langle \phi_i | H_0^2 | \phi_i \rangle = \sum_{i=1}^n \lambda_i \| H_0 | \phi_i \rangle \|^2.$$
(13)

Since

$$\rho^{\alpha}H_{0} = \sum_{i=1}^{n} \lambda_{i}^{\alpha} |\phi_{i}\rangle\langle\phi_{i}|H_{0}$$

and

$$\rho^{1-\alpha}H_0 = \sum_{i=1}^n \lambda_i^{1-\alpha} |\phi_i\rangle \langle \phi_i | H_0,$$

we have

$$\rho^{\alpha}H_{0}\rho^{1-\alpha}H_{0} = \sum_{i,j=1}^{n} \lambda_{i}^{\alpha}\lambda_{j}^{1-\alpha}|\phi_{i}\rangle\langle\phi_{i}|H_{0}|\phi_{j}\rangle\langle\phi_{j}|H_{0}.$$

Thus

$$Tr[\rho^{\alpha}H_{0}\rho^{1-\alpha}H_{0}]$$

$$= \sum_{i,j=1}^{n} \lambda_{i}^{\alpha}\lambda_{j}^{1-\alpha}h_{ij}h_{ji}$$

$$= \sum_{i,j=1}^{n} \lambda_{i}^{\alpha}\lambda_{j}^{1-\alpha}|h_{ij}|^{2}.$$
(14)

By the similar calculations we have

$$Tr[\rho^{\beta}H_{0}\rho^{1-\beta}H_{0}]$$

$$= \sum_{i,j=1}^{n} \lambda_{i}^{\beta}\lambda_{j}^{1-\beta}h_{ij}h_{ji}$$

$$= \sum_{i,j=1}^{n} \lambda_{i}^{\beta}\lambda_{j}^{1-\beta}|h_{ij}|^{2}.$$
(15)

$$Tr[\rho^{\alpha+\beta}H_0\rho^{1-\alpha-\beta}H_0]$$

$$= \sum_{i,j=1}^n \lambda_i^{\alpha+\beta}\lambda_j^{1-\alpha-\beta}h_{ij}h_{ji}$$

$$= \sum_{i,j=1}^n \lambda_i^{\alpha+\beta}\lambda_j^{1-\alpha-\beta}|h_{ij}|^2.$$
(16)

From (5), (13), (14), (15), (16),

$$\begin{split} &I_{\rho,\alpha,\beta}(H) \\ = \frac{1}{2} \sum_{i,j} (\lambda_i + \lambda_i^{\alpha+\beta} \lambda_j^{1-\alpha-\beta} - \lambda_i^{\alpha} \lambda_j^{1-\alpha} - \lambda_i^{\beta} \lambda_j^{1-\beta}) |h_{ij}|^2 \\ &= \frac{1}{2} \sum_i (\lambda_i + \lambda_i - \lambda_i - \lambda_i) |h_{ii}|^2 \\ &+ \frac{1}{2} \sum_{i < j} (\lambda_i + \lambda_i^{\alpha+\beta} \lambda_j^{1-\alpha-\beta} - \lambda_i^{\alpha} \lambda_j^{1-\alpha} - \lambda_i^{\beta} \lambda_j^{1-\beta}) |h_{ij}|^2 \\ &+ \frac{1}{2} \sum_{i < j} (\lambda_j + \lambda_j^{\alpha+\beta} \lambda_i^{1-\alpha-\beta} - \lambda_j^{\alpha} \lambda_i^{1-\alpha} - \lambda_j^{\beta} \lambda_i^{1-\beta}) |h_{ji}|^2 \\ &= \frac{1}{2} \sum_{i < j} (\lambda_i + \lambda_j + f_{\alpha+\beta}(i,j) - f_{\alpha}(i,j) - f_{\beta}(i,j)) |h_{ij}|^2. \end{split}$$

Lemma 3.2:

$$J_{\rho,\alpha,\beta}(H) \ge \sum_{i < j} (\lambda_i + \lambda_j + f_{\alpha+\beta}(i,j) + f_{\alpha}(i,j) + f_{\beta}(i,j)) |h_{ij}|^2.$$

Proof of Lemma 3.2. By (6), (13), (14), (15), (16), we have

$$\begin{split} &J_{\rho,\alpha,\beta}(H)\\ = \quad \frac{1}{2}\sum_{i,j} (\lambda_i + \lambda_i^{\alpha+\beta}\lambda_j^{1-\alpha-\beta} + \lambda_i^{\alpha}\lambda_j^{1-\alpha} + \lambda_i^{\beta}\lambda_j^{1-\beta})|h_{ij}|^2\\ &= \quad \frac{1}{2}\sum_i (\lambda_i + \lambda_i + \lambda_i + \lambda_i)|h_{ii}|^2\\ &+ \frac{1}{2}\sum_{i$$

Lemma 3.3: For any t > 0 and $\alpha, \beta \ge 0, \alpha + \beta \ge 1$ or $\alpha + \beta \le \frac{1}{2}$, the following inequality holds;

$$(t^{1-\alpha-\beta}+1)^2(t^{2\alpha}-1)(t^{2\beta}-1) \ge 16\alpha\beta(t-1)^2.$$
(17)

Proof of Lemma 3.3. It is sufficient to prove (17) for $t \ge 1$ and $\alpha, \beta \ge 0, \alpha + \beta \ge 1$ or $\alpha + \beta \le \frac{1}{2}$. By Lemma 3.3 in [16] we have for $0 \le p \le 1$ and $s \ge 1$,

$$(1-2p)^2(s-1)^2 - (s^p - s^{1-p})^2 \ge 0.$$

Then we can rewrite as follows;

$$(s^{2p}-1)(s^{2(1-p)}-1) \ge 4p(1-p)(s-1)^2.$$

We assume that $\alpha, \beta \ge 0$. We put $p = \alpha/(\alpha + \beta)$ and Since $s^{1/(\alpha+\beta)} = t$. Then

$$(t^{2\alpha} - 1)(t^{2\beta} - 1) \ge \frac{4\alpha\beta}{(\alpha + \beta)^2}(t^{\alpha + \beta} - 1)^2.$$

Then we have

$$(t^{1-\alpha-\beta}+1)^{2}(t^{2\alpha}-1)(t^{2\beta}-1)$$

$$\geq \frac{4\alpha\beta}{(\alpha+\beta)^{2}}(t^{1-\alpha-\beta}+1)^{2}(t^{\alpha+\beta}-1)^{2}.$$
 (18)

We put $\alpha+\beta=k$ and $f(t)=(t^{1-k}+1)(t^k-1)-2k(t-1).$ Then

$$\begin{aligned} f'(t) &= (1-k)t^{-k}(t^k-1) + k(t^{1-k}+1)t^{k-1} - 2k \\ &= (1-k)(1-t^{-k}) + k(1+t^{k-1}) - 2k. \end{aligned}$$

and

$$f''(t) = (1-k)kt^{-k-1} + k(k-1)t^{k-2}$$

= $k(k-1)(t^{k-2} - t^{-k-1}).$

When $k = \alpha + \beta \ge 1$ or $k = \alpha + \beta \le \frac{1}{2}$, it is easy to show that $f''(t) \ge 0$ for $t \ge 1$. Since f'(1) = 0, we have $f'(t) \ge 0$ for $t \ge 1$. And since f(1) = 0, we have $f(t) \ge 0$ for $t \ge 1$. Hence we have for $\alpha + \beta \ge 1$ or $\alpha + \beta \le \frac{1}{2}$,

$$(t^{1-\alpha-\beta}+1)(t^{\alpha+\beta}-1) \ge 2(\alpha+\beta)(t-1).$$

It follows from (18) that we get

$$(t^{1-\alpha-\beta}+1)^2(t^{2\alpha}-1)(t^{2\beta}-1) \ge 16\alpha\beta(t-1)^2.$$

Proof of Theorem 3.1. Since

$$\begin{split} &(t^{1-\alpha-\beta}+1)^2(t^{2\alpha}-1)(t^{2\beta}-1)\\ = &(t+1+t^{\alpha+\beta}+t^{1-\alpha-\beta})^2-(t^\alpha+t^{1-\alpha}+t^\beta+t^{1-\beta})^2, \end{split}$$

we put $t = \frac{\lambda_i}{\lambda_j}$ in (17). Then we have

$$\begin{cases} \frac{\lambda_i}{\lambda_j} + 1 + \left(\frac{\lambda_i}{\lambda_j}\right)^{\alpha+\beta} + \left(\frac{\lambda_i}{\lambda_j}\right)^{1-\alpha-\beta} \end{cases}^2 \\ - \left\{ \left(\frac{\lambda_i}{\lambda_j}\right)^{\alpha} + \left(\frac{\lambda_i}{\lambda_j}\right)^{1-\alpha} + \left(\frac{\lambda_i}{\lambda_j}\right)^{\beta} + \left(\frac{\lambda_i}{\lambda_j}\right)^{1-\beta} \right\}^2 \\ \ge 16\alpha\beta \left(\frac{\lambda_i}{\lambda_j} - 1\right)^2. \end{cases}$$

Then we have

$$\{\lambda_{i} + \lambda_{j} + f_{\alpha+\beta}(i,j) - f_{\alpha}(i,j) - f_{\beta}(i,j)\} \times \{\lambda_{i} + \lambda_{j} + f_{\alpha+\beta}(i,j) + f_{\alpha}(i,j) + f_{\beta}(i,j)\}$$

$$= (\lambda_{i} + \lambda_{j} + f_{\alpha+\beta}(i,j))^{2} - (f_{\alpha}(i,j) + f_{\beta}(i,j))^{2}$$

$$\geq 16\alpha\beta(\lambda_{i} - \lambda_{j})^{2}.$$
(19)

$$Tr[\rho[A, B]] = Tr[\rho[A_0, B_0]]$$

= $2iImTr[\rho A_0 B_0]$
= $2iIm \sum_{i < j} (\lambda_i - \lambda_j) a_{ij} b_{ji}$
= $2i \sum_{i < j} (\lambda_i - \lambda_j) Ima_{ij} b_{ji}$

$$|Tr[\rho[A, B]]| = 2|\sum_{i < j} (\lambda_i - \lambda_j) Ima_{ij} b_{ji}|$$

$$\leq 2\sum_{i < j} |\lambda_i - \lambda_j| |Ima_{ij} b_{ji}|.$$

Then we have

$$|Tr[\rho[A,B]]|^2 \le 4 \left\{ \sum_{i < j} |\lambda_i - \lambda_j| |Ima_{ij}b_{ji}| \right\}^2$$

By (19) and Schwarz inequality,

$$\begin{aligned} &\alpha\beta|Tr[\rho[A,B]]|^2 \\ &\leq 4\alpha\beta\left\{\sum_{i$$

where $K=\lambda_i+\lambda_j+f_{\alpha+\beta}(i,j), L=f_{\alpha}(i,j)+f_{\beta}(i,j).$ Then we have

$$I_{\rho,\alpha,\beta}(A)J_{\rho,\alpha,\beta}(B) \ge \alpha\beta |Tr[\rho[A,B]]|^2.$$

We also have

$$I_{\rho,\alpha,\beta}(B)J_{\rho,\alpha,\beta}(A) \ge \alpha\beta |Tr[\rho[A,B]]|^2.$$

Hence we have the final result (11).

Remark 3.1: We remark that (10) is derived by putting $\beta = 1 - \alpha$ in (11). Then Theorem 3.1 is a generalization of Theorem 2.1 given in [16].

Remark 3.2: When $\alpha, \beta \ge 0$ and $\frac{1}{2} < \alpha + \beta < 1$, we can show an example which Theorem 3.1 does not hold as follows; Let $\alpha = 1/2, \beta = 1/4$ and

$$\rho = \begin{pmatrix} 3/4 & 0 \\ 0 & 1/4 \end{pmatrix}, A = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}, B = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

Then we have

$$U_{\rho,\alpha,\beta}(A)U_{\rho,\alpha,\beta}(B) = 0.004487,$$

 $\alpha\beta|Tr[\rho[A,B]]|^2 = 0.125.$

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